

Scalar Quantum Fields in General Relativity

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Ref: Mukhanov, V.F., and Winitzki, S., *Introduction to Quantum Effects in Gravity*, Cambridge (2007), Chap. 5

	<u>Minimally Coupled</u>		<u>Non-Minimally Coupled Example</u>		<u>Comment</u>
	<u>Real Scalar</u>	<u>Gravity</u>	<u>Real Scalar</u>	<u>Gravity</u>	
Field	ϕ	$g_{\mu\nu}$	ϕ	$g_{\mu\nu}$	
Back-ground	Physical length (measured in meters with meter sticks) in x_1 direction is $dx_1 = \sqrt{-g_{11}} dx_1$ Coordinate length = dx_1 (The length of this distance in meters measured with meter sticks depends on the particular generalized coordinate grid chosen, i.e., on g_{11} , which represents that grid in the x_1 direction.) Physical 4D volume in an orthogonal (i.e., $g_{\mu\nu}$ is diagonal) spacetime coordinate system is $dV = dx_1 dx_2 dx_3 dx_0$ $= \sqrt{-g_{11}} dx_1 \sqrt{-g_{22}} dx_2 \sqrt{-g_{33}} dx_3 \sqrt{g_{00}} dx_0 = \sqrt{-g} dx_1 dx_2 dx_3 dx_0 = \sqrt{-g} dV = \sqrt{-g} d^4 x$			g_{11} is negative; phys value has “~” underneath g is Det $g_{\mu\nu}$; dV is coord volume	
Procedure	In standard QFT 1) $\eta_{\mu\nu} \rightarrow g_{\mu\nu}$ 2) “ , ” \rightarrow “ ; ” 3) $d^4 x \rightarrow \sqrt{-g} d^4 x$	Choose \mathcal{L}_{grav} as in classical GR	In this example, as in 2 nd column, but add extra term.	In this example, \mathcal{L}_{grav} of 3 rd column.	In QFT we change to natural units, where $c = \hbar = 1$
Action S	$S_m = \int \underbrace{\mathcal{L}_m}_{\text{phys Lagr}} \underbrace{\sqrt{-g} d^4 x}_{\text{phys } dV}$ $= \int \underbrace{\mathcal{L}_m}_{\text{coord coord Lagr}} \underbrace{d^4 x}_{dV}$	$S_{grav} = \int \underbrace{\mathcal{L}_{grav}}_{\text{phys Lagr}} \underbrace{\sqrt{-g} d^4 x}_{\text{phys } dV}$ $= \int \mathcal{L}_{grav} d^4 x$	As in 2 nd column	As in 3 rd column	$\sqrt{-g}$ of 3) above included in \mathcal{L}_m and \mathcal{L}_{grav}
Lagrangian \mathcal{L}	$\mathcal{L}_m = \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} \phi_{,\mu} \phi_{,\nu} - V(\phi) \right)$	$\mathcal{L}_{grav} = -\frac{\sqrt{-g}}{8\pi G} (R + 2\Lambda)$	$\mathcal{L}_m = \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} \phi_{,\mu} \phi_{,\nu} - V(\phi) - \frac{\xi}{2} R \phi^2 \right)$	In this example, \mathcal{L}_{grav} of 3 rd column	Other non-min \mathcal{L}_m , \mathcal{L}_{grav} could have other terms in $R_{\alpha\beta\gamma\delta}$
Free field V	$V(\phi) = \frac{1}{2} m^2 \phi^2$		As in 2 nd column		“Free”, but ϕ , grav coupled
Total \mathcal{L}	$\mathcal{L} = \mathcal{L}_m + \mathcal{L}_{grav}$	$\mathcal{L} = \mathcal{L}_m + \mathcal{L}_{grav}$	As in 2 nd column	As in 3 rd column	
Variation of S	From $\delta S = 0$, where \mathcal{L}_{grav} has terms in $g_{\alpha\beta,\mu\nu}$ due to R , one gets the Euler-Lagrange eqs below. Gravity eq has an extra term beyond the more familiar Euler-Lagrange eq due to 2 nd derivative terms in \mathcal{L}_{grav} . Note the math derivation is based on the integrand used with $d^4 x$, so it uses $\mathcal{L} = \mathcal{L} \sqrt{-g}$.				
Euler-Lagrange equation	$\partial_\mu \frac{\partial \mathcal{L}}{\partial \phi_{,\mu}} - \frac{\partial \mathcal{L}}{\partial \phi} = 0$	$-\partial_\mu \partial_\nu \frac{\partial \mathcal{L}}{\partial g_{\alpha\beta,\mu\nu}}$ $+ \partial_\mu \frac{\partial \mathcal{L}}{\partial g_{\alpha\beta,\mu}} - \frac{\partial \mathcal{L}}{\partial g_{\alpha\beta}} = 0$	As in 2 nd column	As in 3 rd column	From $\delta S = 0$
Note:	Derivation to get gravity result below is laborious. See Mukhanov & Minitzki pgs. 229-232.				

Equation of motion	$(\sqrt{-g} g^{\alpha\beta} \phi_{,\beta})_{,\alpha} + \sqrt{-g} \frac{\partial V}{\partial \phi} = 0$	$G_{\alpha\beta} + g_{\alpha\beta} \Lambda = R_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} R + g_{\alpha\beta} \Lambda = 8\pi G T_{\alpha\beta}$	$(\sqrt{-g} g^{\alpha\beta} \phi_{,\beta})_{,\alpha} + \sqrt{-g} \left(\frac{\partial V}{\partial \phi} + \xi R \phi \right) = 0$	As in 3 rd column w different $T_{\alpha\beta}$	From Euler-Lagrange equation
Stress-energy part of above		$T_{\alpha\beta} = \frac{2}{\sqrt{-g}} \frac{\partial \mathcal{L}_m}{\partial g^{\alpha\beta}} = \phi_{,\alpha} \phi_{,\beta} - g_{\alpha\beta} \left(\frac{1}{2} g^{\mu\nu} \phi_{,\mu} \phi_{,\nu} - V(\phi) \right)$		$T_{\alpha\beta} = \frac{2}{\sqrt{-g}} \times \left(\begin{array}{l} \partial_\mu \partial_\nu \frac{\partial \mathcal{L}_m}{\partial g_{\alpha\beta, \mu\nu}} \\ - \partial_\mu \frac{\partial \mathcal{L}_m}{\partial g^{\alpha\beta, \mu}} \\ + \frac{\partial \mathcal{L}_m}{\partial g^{\alpha\beta}} \end{array} \right)$	$\leftarrow \mathcal{L}_m$ in 5 th column from "Lagrangian \mathcal{L} " row above, 4 th column
Covariant form of eq of motion	$\phi_{;\alpha}^{\alpha} + \frac{\partial V}{\partial \phi} = 0$	$G_{\alpha\beta} + g_{\alpha\beta} \Lambda = R_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} R + g_{\alpha\beta} \Lambda = 8\pi G T_{\alpha\beta}$	$\phi_{;\alpha}^{\alpha} + \frac{\partial V}{\partial \phi} + \xi R \phi = 0$	As in 3 rd column w different $T_{\alpha\beta}$	Gravity field eq of motion above already covariant
Above are	GR version of Klein-Gordon eq	Einstein's field equation			
<u>E/m Interactions with (Complex) Scalar Field</u>					
Procedure	Initially as above with free field, but now need complex scalar field, as real scalar field is charge neutral.				
Action S	$S_m = \int \mathcal{L}_m d^4x$	All gravity as above in free scalar field			
Free field Lagrangian \mathcal{L}	$\mathcal{L}_m = \sqrt{-g} \times \left(\begin{array}{l} \frac{1}{2} g^{\mu\nu} \phi_{,\mu}^* \phi_{,\nu} \\ - \frac{1}{2} m^2 \phi^* \phi \end{array} \right)$				For free field $V(\phi^*, \phi) = \frac{1}{2} m^2 \phi^* \phi$
Interaction \mathcal{L}_m	with $\phi_{,\mu} \rightarrow D_\mu \phi$ $\mathcal{L}_m = \sqrt{-g} \times \left(\begin{array}{l} \frac{1}{2} g^{\mu\nu} (D_\mu \phi)^* D_\nu \phi \\ - \frac{1}{2} m^2 \phi^* \phi \end{array} \right)$				$D_\mu = \partial_\mu + iA_\mu$
Re-write interact \mathcal{L}_m	$\mathcal{L}_m = \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} \phi_{,\mu}^* \phi_{,\nu} - \frac{1}{2} m^2 \phi^* \phi - \frac{1}{2} g^{\alpha\beta} i A_\alpha (\phi^* \phi_{,\beta} - \phi \phi_{,\beta}^*) + \frac{1}{2} g^{\alpha\beta} A_\alpha A_\beta \phi^* \phi \right)$				
Interaction $\mathcal{L}_{e/m}$	$F_{\mu\nu} = A_{\nu,\mu} - A_{\mu,\nu} = A_{\nu,\mu} - A_{\mu,\nu}$ Standard QFT $\mathcal{L}_{e/m} = -\frac{1}{16\pi} \eta^{\alpha\beta} \eta^{\mu\nu} F_{\alpha\mu} F_{\beta\nu} \rightarrow -\frac{1}{16\pi} g^{\alpha\beta} g^{\mu\nu} F_{\alpha\mu} F_{\beta\nu}$				$\Gamma_{\beta\gamma}^\alpha$ drop out of $F_{\mu\nu}$
Total \mathcal{L}	$\mathcal{L} = \mathcal{L}_m + \mathcal{L}_{grav} + \mathcal{L}_{e/m}$				
Equations of Motion	Use above \mathcal{L} in Euler-Lagrange eq for 3 fields, ϕ , $g_{\mu\nu}$, and A_μ . Get 3 coupled eqs.				