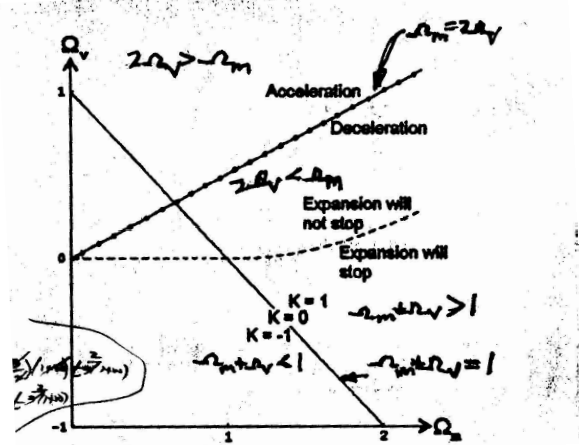


Brief Summary of Cosmology

See Jordan, *Am J Phys* **73**(7) 653-662 (July 2005) for much of this. by R. Klauber

<u>Central Point</u>	<u>Derivation/Remark</u>
<p>Conservation of Energy: $d(\rho V) = -pdV$</p> <p>Energy Density: $\dot{\rho} = -3(\rho + p)\dot{a}/a = -3(\rho + p)H$</p> <p>$\rho(a)$ w/ $p=w\rho$: $\rho_w \propto 1/a^{3(1+w)}$</p> <p>Rad $\rho_{w=1/3} \propto 1/a^4$ Matter $\rho_{w=0} \propto 1/a^3$ Vac $\rho_{w=-1} \propto \text{const}$</p>	<p>3 ways to Energy Density relation:</p> <ol style="list-style-type: none"> 1) Divide top eq by dt, use $V=ka^3$ 2) Freid metric in $T^{\mu\nu}_{;\nu} = 0$ for $\mu = 0$. 3) $T^{\hat{0}\hat{0}}$ & $T^{\hat{x}\hat{x}}$ field eqs combined. <p>$\rho(a)$: ρ_w into top eq, w/ $V=ka^3$ & $p=w\rho$ $c^2\rho = \rho$ in all this ($c = 1$)</p>
<p>Basic Eqs: Freidman Universe (homogeneous, isotropic)</p> <p>(A) $\left(\frac{\dot{a}}{a}\right)^2 = H^2 = \frac{8\pi G\rho}{3} + \frac{\Lambda}{3} - k\left(\frac{a_0}{a}\right)^2$ Freidman (initial val) Eq.</p> <p>(B) $\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{3}$ Dynamic eq.</p>	<p>Can rep Λ in ρ_{vac} and take $\Lambda = 0$ here.</p> <p>Robertson-Walker metric. Einstein eq for T^{00}</p> <p>Einstein eq for T^{xx} plus (A), or energy density eq plus (A)</p>
<p>Radiation: $\lambda \propto a$ $\bar{E}_\gamma = 2.7kT \propto 1/\lambda \propto 1/a$</p> <p>$n_\gamma \propto 1/a^3$ $\rho_\gamma = n_\gamma \bar{E}_\gamma = 2.7kn_\gamma T \propto 1/a^4$ also $p = \rho_\gamma/3$</p>	
<p>$k = 0$ Time Dependence ($\Lambda = 0$):</p> <p>Rad: $a \propto t^{1/2}$ Matter: $a \propto t^{2/3}$</p> <p>Vacuum: $a \propto e^{Ht}$ General: $a \propto t^{2/3(1+w)}$ (for $\rho_v = -p = \text{const}$)</p>	<p>(A) $\rightarrow \dot{a} = \sqrt{\frac{8\pi G\rho}{3}}a = Ha$ and use various values for $\rho(a)$ above.</p>
<p>Critical Densty: $\rho_c = \frac{3H^2}{8\pi G}$ $\Omega = \rho / \rho_c = \Omega_\gamma + \Omega_m + \Omega_v$ $\Omega_i = \rho_i / \rho$</p> <p>Rad (early univ) $\rho \propto 1/a^4$, $1 - \frac{1}{\Omega} \propto a^2$ and $\rho > \rho_c$ grows fast</p> <p>For inflation $\rho = \text{const}$, $1 - \frac{1}{\Omega} \propto \frac{k}{a^2}$ and $\rho \rightarrow \rho_c$ fast</p>	<p>From (A), $\Omega = 1$, if $k = 0$.</p> <p>(A) w/ k and $\Lambda = 0$, rearranged gives this relation in $1/\Omega$.</p> <p>If $k/a^2 \rightarrow 0$, then $\Omega \rightarrow 1$. If $\Omega=1$ at any time, then $k=0$ and $\Omega=1$ for all time</p>
<p>Acceleration: $\ddot{a} = -\frac{4\pi G}{3}(\rho_m - 2\rho_v)a$ $\ddot{a} > 0$ if $2\rho_v > \rho_m$</p>	<p>(B) for mass and vacuum, assuming vac like cosmol const with $w = -1$.</p> <p>More generally, in (B), if $p < -\rho/3$ (i.e., $w < -1/3$), where p and ρ are total values, get acceleration.</p>



<p>Travel Time: <u>Basic principles:</u> 1) $\Delta t = \int_a^{a_0} \frac{dt(a')}{da'} da'$</p> <p>2) Find da/dt from (B) and invert</p> <p><u>Subtleties:</u> $r = \frac{a}{a_0} \leq 1$ $\tau = \frac{t}{1/H_0} = H_0 t$ $r = \frac{\lambda}{\lambda_0} = \frac{1}{1+z}$</p> <p><u>Result:</u> $\Delta \tau \left(r = \frac{1}{1+z} \right) = \int_r^1 \frac{d\tau(r')}{dr'} dr'$</p> <p>where $\frac{d\tau}{dr} = \left(\frac{\Omega_{m_0}}{r} + \Omega_{v_0} r^2 + 1 - \Omega_{m_0} - \Omega_{v_0} \right)^{-1/2}$</p> <p>for vac with $w \neq -1$, get diff relation.</p>	<p>Ref Jordan.1) Looks flawed to me. Easier to start w basic princ in Hubble plot box</p> <p>Integ (B). Here, result assumes $w_v = -1$.</p> <p>For integ const, need to evaluate ka_0^2 term via Hubble relation now, i.e.,</p> $\frac{da}{dt} = H_0 a_0 \rightarrow \frac{d(a/a_0)}{H_0 dt} = 1 = \frac{dr}{d\tau} \text{ now}$ <p>This gives constant terms in last line in LHS box.</p>
<p>Hubble Plots: <u>Def:</u> d_l = luminosity dist = distance to source if flat, static universe for the luminosity we see for standard candle.</p> <p><u>Basic principle:</u> luminosity dist = $\int_r^1 \frac{cd\tau(r')}{dr'} dr'$ (see above)</p> <p><u>Corrections:</u></p> <p>1) expansion on dist $1/r'$ factor $\chi = \int_r^1 \frac{1}{r'} \frac{cd\tau(r')}{dr'} dr'$</p> <p>2) curved space: $k=0$, $\sigma = \chi$; $k=1$, $r_c \sin \frac{\chi}{r_c}$; $k=-1$, $r_c \sinh \frac{\chi}{r_c}$</p> <p>3) expansion on photons: a) decrease intensity per γ by $\lambda/\lambda_0 = r$ b) decrease γ/sec arriving by $r = \lambda/\lambda_0$</p> <p><u>Thus:</u> intensity $\propto 1/[4\pi(d_l)^2] \propto 1/[4\pi(\sigma)^2(1+z)^2]$</p> <p><u>Subtlety:</u> d_l defined in terms of units of H_0, $d_l = (1+z) \frac{\sigma}{H_0}$</p> <p><u>Finally:</u> intensity $\propto \frac{1}{4\pi \left(\frac{\sigma}{H_0} \right)^2 (1+z)^2}$</p>	<p>Original Hubble plots, vel vs dist. Now, d_l vs z</p> <p>= 13.7 bill lt-yrs for pg 6 example</p> <p>= 42 bill lt-yrs for pg 6 example</p> $r_c = \frac{1}{\sqrt{1-\Omega_0}}$ <p><u>Value:</u> d_l vs z curve depends on Ω_{m_0} and Ω_{v_0} (for $w_v = -1$ case), so data helps determine them and from (B), \ddot{a}. For $w_v \neq -1$, will get diff $d\tau/dr$, so diff σ and diff curve.</p> <p>Supernovae (of known intrinsic brightness) shows accel and $w_v = -1$ const within 10% over most of hist.</p>